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## Positive solutions of nonlinear elliptic equations in the case of critical Sobolev exponent

We report on a recent work with L. Nirenberg [5]. Consider the following problem. Assume  $\Omega \subset \mathbb{R}^N$ ,  $N \geq 3$ , is a bounded (smooth) domain. Find a (smooth) function  $u$  satisfying

$$\begin{cases} u > 0 & \text{in } \Omega \\ -\Delta u = u^p + f(u) & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega, \end{cases} \quad (1)$$

where  $p = \frac{N+2}{N-2}$  and  $f(u)$  is a "lower order perturbation" with  $f(0) = 0$ ; a typical example is  $f(u) = \lambda u$  ( $\lambda \in \mathbb{R}$ ). The exponent  $p = \frac{N+2}{N-2}$  is critical from the point of view of the variational formulation. Indeed, solutions of (1) correspond to critical points of the functional  $\frac{1}{2} \int |\nabla u|^2 - \frac{1}{p+1} \int u^{p+1} - \int F(u)$  where  $F$  is a primitive of  $f$  and  $p+1 = \frac{2N}{N-2}$  is the Sobolev exponent for the embedding  $H_0^1(\Omega) \subset L^{p+1}(\Omega)$ .

Our lecture is divided as follows. First we recall some results concerning the *easy* case where  $p < \frac{N+2}{N-2}$  and  $f(u) = \lambda u$ . Then, we consider the case where  $p = \frac{N+2}{N-2}$  and  $f(u) = \lambda u$ . Finally we turn to the case where  $f$  is nonlinear.

Our interest in problem (1) comes from the fact that it presents some similarities with the Yamabe problem in geometry; see e.g. Trudinger [11] and Th. Aubin [2].

1. THE CASE  $p < \frac{N+2}{N-2}$ .

Throughout Section 1 we assume that  $p < \frac{N+2}{N-2}$ . Clearly, there is a solution  $u$  of

$$\begin{cases} u > 0 & \text{in } \Omega \\ -\Delta u = u^p & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases} \quad (2)$$

Indeed, consider the following minimization problem

$$\inf_{v \in H_0^1} \left\{ \frac{\|\nabla v\|_{L^2}^2}{\|v\|_{L^{p+1}}^2} \right\} \quad (3)$$

Since the injection  $H_0^1 \subset L^{p+1}$  is compact, the infimum in (3) is achieved by some  $v_0$ . We may always assume that  $v_0 \geq 0$  (otherwise replace  $v_0$  by  $|v_0|$ ) and that  $\|v_0\|_{L^{p+1}} = 1$ . Thus we obtain a Lagrange multiplier  $\mu \in \mathbb{R}$  such that

$$-\Delta v_0 = \mu v_0^p \quad (4)$$

and  $\mu = \int |\nabla v_0|^2 > 0$ .

By stretching  $v_0$  we find a function  $u$  satisfying

$$\begin{cases} u \geq 0 & \text{on } \Omega, & u \not\equiv 0 \\ -\Delta u = u^p & \text{on } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases} \quad (5)$$

[more precisely  $u = kv_0$  satisfies (5) provided  $k = \mu^{\frac{1}{p-1}}$ ]. It follows from the strong maximum principle that  $u > 0$  in  $\Omega$ .

(The question of uniqueness for problem (2) is open when  $\Omega$  is starshaped.

When  $\Omega$  is an annulus and  $p$  is close to  $\frac{N+2}{N-2}$  the solution of (2) need not be unique ; in fact there exist both spherical and non-spherical solutions, see [5]).

Let  $\lambda_1$  denote the first eigenvalue of  $-\Delta$  with zero Dirichlet boundary condition. Consider now the following problem : find  $u$  such that

$$\begin{cases} u > 0 & \text{in } \Omega \\ -\Delta u = u^p + \lambda u & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases} \quad (6)$$

Then for each  $\lambda \in (-\infty, \lambda_1)$  there is a solution of (6). Indeed

$$\inf_{v \in H_0^1} \left\{ \frac{\|\nabla v\|_{L^2}^2 - \lambda \|v\|_{L^2}^2}{\|v\|_{L^{p+1}}^2} \right\} \quad (7)$$

is achieved by some  $v_0$  satisfying  $v_0 \geq 0$  on  $\Omega$  and  $\|v_0\|_{L^{p+1}} = 1$ . Moreover there is a Lagrange multiplier  $\mu \in \mathbb{R}$  such that

$$-\Delta v_0 - \lambda v_0 = \mu v_0^p.$$

Thus  $\mu = \int |\nabla v_0|^2 - \lambda \int v_0^2 > 0$  (since  $\lambda < \lambda_1$ ). By stretching  $v_0$  as above we obtain a solution of (6).

The restriction  $\lambda \in (-\infty, \lambda_1)$  is essential. Indeed suppose  $u$  is a solution of (6). Let  $\phi_1 > 0$  in  $\Omega$  be an eigenfunction corresponding to  $\lambda_1$ . We have

$$\lambda_1 \int u \phi_1 = \int u^p \phi_1 + \lambda \int u \phi_1 > \lambda \int u \phi_1$$

and thus  $\lambda < \lambda_1$ .

## 2. THE CASE $p = \frac{N+2}{N-2}$ AND $f(u) = \lambda u$

Throughout Sections 2 and 3 we assume that  $p = \frac{N+2}{N-2}$ . We consider now the following problem : find  $u$  such that

$$\begin{cases} u > 0 & \text{in } \Omega \\ -\Delta u = u^p + \lambda u & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases} \quad (\lambda \in \mathbb{R}) \quad (8)$$

The argument we have used on Section 1 does not hold anymore since the injection  $H_0^1 \subset L^{p+1}$  is *not* compact. In fact we know, by a result of Pohozaev [9], that if  $\Omega$  is *starshaped* and  $\lambda = 0$  there is *no* solution of (8). Using the same argument as in Pohozaev [9] one proves :

Theorem 0 : Assume  $\Omega$  is *starshaped* and  $\lambda \leq 0$ . Then there is no solution of (8).

Remark 1 : On the other hand if  $\Omega$  is an *annulus* then for every  $\lambda \in (-\infty, \lambda_1)$  there is a spherical solution of (8). Indeed consider

$$\inf_{v \in H_r^1} \left\{ \frac{\|\nabla v\|_{L^2}^2 - \lambda \|v\|_{L^2}^2}{\|v\|_{L^{p+1}}^2} \right\} \quad (9)$$

where  $H_r^1 = \{v \in H_0^1(\Omega) \text{ and } v \text{ is spherically symmetric}\}$ .

The infimum in (9) is achieved since the injection of  $H_r^1$  into  $L^{p+1}$  is *compact*. Thus, after stretching, we obtain a spherical solution of (8).

Our main results are the following

Theorem 1 : Assume  $\Omega \subset \mathbb{R}^N$ ,  $N \geq 4$ , is any (smooth) bounded domain. Then for every  $\lambda \in (0, \lambda_1)$  there exists a solution of (8). Moreover

$$\inf_{v \in H_0^1} \left\{ \frac{\|\nabla v\|_{L^2}^2 - \lambda \|v\|_{L^2}^2}{\|v\|_{L^{p+1}}^2} \right\} \quad (10)$$

is achieved.

Theorem 2 : Assume  $\Omega$  is a ball in  $\mathbb{R}^3$ . Then for every  $\lambda \in (\frac{\lambda_1}{4}, \lambda_1)$  there is a solution of (8) ; moreover the infimum in (10) is achieved. When  $\lambda \leq \frac{\lambda_1}{4}$  there is no solution of (8).

Remark 2 : The difference between dimension  $N = 3$  and dimension  $N \geq 4$  is quite suprising. We have no simple explanation for it.

Remark 3 : When  $N \geq 3$  and  $\lambda \geq \lambda_1$  there is no solution of (8) (see Section 1).

When  $N \geq 3$  and  $\Omega$  is *starshaped* there is no solution of (8) for  $\lambda \leq 0$  (by Theorem 0).

Remark 4 : The generalization of Theorem 2 for starshaped domains is not known.

Before we sketch the proofs we present some facts about Sobolev spaces :

- a) Define the best Sobolev constant  $S$  to be

$$S = \inf_{v \in H_0^1} \left\{ \frac{\|\nabla v\|_{L^2}^2}{\|v\|_{L^{p+1}}^2} \right\} \quad (11)$$

In principle  $S$  depends on  $\Omega$ ; but in fact  $S$  depends only on  $N$ . This is an easy consequence of the invariance under scaling of the ratio  $\frac{\|\nabla v\|_{L^2}^2}{\|v\|_{L^{p+1}}^2}$  (that is, the ratio is unchanged if we replace  $u(x)$  by  $u_k(x) = u(kx)$ ).

b) The infimum in (11) is never achieved, on any bounded domain. Indeed, suppose that the infimum in (11) is achieved by some function  $v_0 \geq 0$ . Let  $\tilde{\Omega}$  be a ball containing  $\Omega$  and set

$$\tilde{v}_0 = \begin{cases} v_0 & \text{in } \Omega \\ 0 & \text{in } \tilde{\Omega} \setminus \Omega \end{cases}$$

Thus, for  $\tilde{\Omega}$  the infimum in (11) is achieved at  $v_0$  and we find

$$\begin{cases} -\Delta \tilde{v}_0 = \mu v_0^p & \text{on } \Omega \\ v_0 = 0 & \text{on } \partial\Omega \end{cases}$$

for some constant  $\mu > 0$ . This contradicts Pohozaev's Theorem.

c) When  $\Omega = \mathbb{R}^N$  the infimum in (11) is achieved by the function  $u(x) = \frac{1}{(1+|x|^2)^{\frac{N-2}{2}}}$  or - after scaling - by any of the functions

$$u(x) = \frac{1}{(\varepsilon + |x|^2)^{\frac{N-2}{2}}} \quad (\varepsilon > 0).$$

see Aubin [3], Talenti [10], Lieb [8].

The following Lemma plays a crucial role in the proof of Theorem 1.

Lemma 1 : Assume  $N \geq 4$ . Then, for every  $\lambda \in (0, \lambda_1)$  we have

$$S_\lambda \equiv \inf_{v \in H_0^1} \left\{ \frac{\|\nabla v\|_{L^2}^2 - \lambda \|v\|_{L^2}^2}{\|v\|_{L^{p+1}}^2} \right\} < S \quad (12)$$

The proof of Lemma 1 is rather technical ; for details see [5]. The main idea - borrowed from Aubin [1] - consists of estimating the ratio

$$Q(u) = \frac{\|\nabla u\|_{L^2}^2 - \lambda \|u\|_{L^2}^2}{\|u\|_{L^{p+1}}^2}$$

for  $u(x) = u_\epsilon(x) = \frac{\phi(x)}{(\epsilon + |x|^2)^{\frac{N-2}{2}}}$  where  $\phi \in \mathcal{D}_+(\Omega)$  is a fixed function such

that  $\phi(x) \equiv 1$  near 0 (assuming  $0 \in \Omega$ ).

A straightforward computation gives the following expansion as  $\epsilon \rightarrow 0$  :

$$Q(u_\epsilon) = S + O(\epsilon^{\frac{N}{2}-1}) - \lambda C \epsilon \quad \text{when } N \geq 5$$

$$Q(u_\epsilon) = S + O(\epsilon) - \lambda C \epsilon |\log \epsilon| \quad \text{when } N = 4$$

where  $C > 0$  is a constant. In both cases we see that  $Q(u_\epsilon) < S$  for  $\epsilon > 0$  sufficiently small. We shall also use the following measure theoretic lemma.

Lemma 2 : (Brezis-Lieb [4]) Suppose  $(v_j)$  is a sequence in  $L^q(\Omega)$  with  $1 \leq q < \infty$  such that  $\|v_j\|_{L^q}$  remains bounded and  $v_j(x) \rightarrow v(x)$  a.e. on  $\Omega$ . Then

$$\lim_{j \rightarrow \infty} \left\{ \int |v_j|^q - \int |v_j - v|^q \right\} = \int |v|^q \quad (13)$$

Proof of Theorem 1 : Choose a minimizing sequence  $(v_j)$  for (10) such that

$$v_j \geq 0 \text{ on } \Omega, \quad \|v_j\|_{L^{p+1}} = 1 \quad (14)$$

$$\int |v_j|^2 - \lambda \int v_j^2 = S_\lambda + o(1) \quad (15)$$

Since  $v_j$  is bounded in  $H_0^1$  we may assume, for a subsequence, that

$$v_j \rightharpoonup v \text{ weakly in } H_0^1$$

$$v_j \rightarrow v \text{ a.e. on } \Omega$$

$$v_j \rightarrow v \text{ strongly in } L^2$$

Set  $w_j = v_j - v$  so that  $w_j \rightharpoonup 0$  weakly in  $H_0^1$ .

By Lemma 2 (applied with  $q = p+1$ ) we have

$$\lim \|w_j\|_{L^{p+1}}^{p+1} = 1 - \|v\|_{L^{p+1}}^{p+1}.$$

Thus

$$\begin{aligned} 1 &= (\|v\|_{L^{p+1}}^{p+1} + \|w_j\|_{L^{p+1}}^{p+1})^{\frac{2}{p+1}} + o(1) \leq \\ &\leq \|v^2\|_{L^{p+1}} + \|w_j\|_{L^{p+1}}^2 + o(1) \end{aligned} \quad (16)$$



On the other hand (since  $w_j \rightarrow 0$  weakly in  $H_0^1$ ) we have

$$\int |\nabla v_j|^2 = \int |\nabla v|^2 + \int |\nabla w_j|^2 + o(1) \quad (17)$$

Combining (15) (16) and (17) we obtain

$$\int |\nabla v|^2 - \lambda \int v^2 + \int |\nabla w_j|^2 \leq S_\lambda [\|v\|_{L^{p+1}}^2 + \|w_j\|_{L^{p+1}}^2] + o(1) \quad (18)$$

By definition of  $S_\lambda$  we have  $\int |\nabla v|^2 - \lambda \int v^2 \geq S_\lambda \|v\|_{L^{p+1}}^2$  and therefore

$$\begin{aligned} \int |\nabla w_j|^2 &\leq S_\lambda \|w_j\|_{L^{p+1}}^2 + o(1) \\ &\leq \frac{S_\lambda}{S} \int |\nabla w_j|^2 + o(1). \end{aligned}$$

Since  $S_\lambda < S$  (by Lemma 1), it follows that  $\int |\nabla w_j|^2 \rightarrow 0$ . Consequently  $v_j \rightarrow v$  strongly in  $H_0^1$  (and in  $L^{p+1}$ ). Passing to the limit in (14) and (15) we conclude that the infimum in (11) is achieved by  $v$ . After stretching we obtain a solution of (8).

In the proof of Theorem 2 we use

Lemma 3 : Assume  $\Omega = \{x \in \mathbb{R}^3 ; |x| < 1\}$ . Then for each  $\lambda \in (\frac{\lambda_1}{4}, \lambda_1)$  we have

$$S_\lambda \equiv \inf_{v \in H_0^1} \left\{ \frac{\|\nabla v\|_{L^2}^2 - \lambda \|v\|_{L^2}^2}{\|v\|_{L^6}^2} \right\} < S \quad (19)$$



Proof : We estimate the ratio

$$Q(u) = \frac{\| \nabla u \|_{L^2}^2 - \lambda \| u \|_{L^2}^2}{\| u \|_{L^6}^2}$$

$$\text{for } u(x) = u_\varepsilon(x) = \frac{\cos(\frac{\pi}{2}|x|)}{(\varepsilon + |x|^2)^{1/2}}.$$

A technical computation (see [5]) gives the following expansion as  $\varepsilon \rightarrow 0$  :

$$Q(u_\varepsilon) = S + C\sqrt{\varepsilon} \left( \frac{\pi^2}{4} - \lambda \right) + o(\sqrt{\varepsilon})$$

where  $C > 0$  is a constant. Therefore if  $\lambda > \frac{\lambda_1}{4}$  (note that here  $\lambda_1 = \pi^2$ ) we see that  $Q(u_\varepsilon) < S$  for  $\varepsilon > 0$  sufficiently small.

Proof of Theorem 2 : The same argument as in the proof of Theorem 1 shows that for every  $\lambda \in (\frac{\lambda_1}{4}, \lambda_1)$  the infimum in (19) is achieved. Consequently we obtain a solution of (8) for each  $\lambda \in (\frac{\lambda_1}{4}, \lambda_1)$ . Next we must show that no solution of (8) exists for  $\lambda \geq \frac{\lambda_1}{4}$ . By a result of Gidas-Nirenberg [7] we know that any solution  $u$  of (8) in a ball must be spherically symmetric. We write  $u(x) = u(r)$  ( $r = |x|$ ) and so we have

$$-u'' - \frac{2}{r}u' = u^5 + \lambda u \quad \text{on } (0,1) \quad (20)$$

$$u'(0) = u(1) = 0. \quad (21)$$

Then we use an argument "à la Pohozaev" but with more complicated multipliers. Namely we multiply (20) through by

$$r^2[r \cos \pi r - b \sin \pi r]u'$$

and then by

$$r[-\frac{r}{2}(1+b\pi)\cos \pi r - \frac{\pi r^2}{2} \sin \pi r + b \sin \pi r] u$$

for some appropriate constant  $b$ . Integrating by parts and combining the two equalities leads to  $\lambda > \frac{\pi^2}{4}$ ; for more details see [5].

### 3. THE GENERAL CASE, $-\Delta u = u^p + f(u)$ WITH $p = \frac{N+2}{N-2}$ .

Here again we take  $p = \frac{N+2}{N-2}$ . Assume  $f$  is a  $C^1$  function on  $[0, +\infty)$  such that

$$f(0) = 0, \quad f(u) \geq 0 \quad \forall u \geq 0 \quad (22)$$

$$\lim_{u \rightarrow \infty} \frac{f(u)}{u^p} = 0 \quad (23)$$

$$f'(0) < \lambda_1 \quad (24)$$

The problem is to find a function  $u$  satisfying

$$\begin{cases} u > 0 & \text{on } \Omega \\ -\Delta u = u^p + f(u) & \text{on } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases} \quad (25)$$

Our main results are the following

Theorem 3 : Assume  $N \geq 5$ , (22), (23), (24) and

$$f \neq 0 \quad (26)$$

Then there is a solution of (25).

Theorem 4 : Assume  $N = 4$ , (22), (23), (24) and either

$$f'(0) > 0 \quad (27)$$

or

$$\liminf_{u \rightarrow +\infty} \frac{f(u)}{u} > 0 \quad (28)$$

Then there is a solution of (25).

Theorem 5 : Assume  $N = 3$ , (22), (23), (24)

$$\lim_{u \rightarrow +\infty} \frac{f(u)}{u^3} = +\infty$$

Then there is a solution of (25).

Remark 5 : Theorem 3, 4 and 5 admit appropriate extensions to the case where  $f$  depends also on  $x$  with  $f(x,0) \equiv 0$ . They may be used in order to prove the following :

Theorem 6 : Assume  $N \geq 3$ . Then there is a constant  $\lambda^* > 0$  such that the problem

$$\begin{cases} u > 0 & \text{in } \Omega \\ -\Delta u = \lambda(1+u)^p & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases} \quad p = \frac{N+2}{N-2} \quad (29)$$

has at least two solutions for each  $\lambda \in (0, \lambda^*)$  and no solution for  $\lambda > \lambda^*$ .

(A similar result for  $p < \frac{N+2}{N-2}$  had been obtained earlier by Crandall-Rabinowitz [6]).

The idea of the proof is the following. Firstly one obtains (easily) a minimal solution  $\underline{u}$  of (29) for every  $\lambda \in (0, \lambda^*)$  (see e.g. [6]). Then one looks for a second solution of (29) of the form  $u = \underline{u} + v$  with  $v > 0$  on  $\Omega$ . Thus  $v$  satisfies  $-\Delta v = \lambda(1+\underline{u}+v)^p - \lambda(1+\underline{u})^p \equiv \lambda v^p + f(x, v)$  and we are reduced to a problem of the type (25).

The proofs of Theorems 3, 4 and 5 involve two ingredients. Firstly a geometrical result which is a variant of the Ambrosetti-Rabinowitz [1] mountain pass Lemma without the (PS) condition (see Lemma 4). Secondly a technical Lemma which has the same flavour as Lemma 1 or Lemma 3 (see Lemma 5).

Lemma 4 : Assume  $\Phi$  is a  $C^1$  function on a Banach space  $E$  such that

$$\Phi(0) = 0 \quad (30)$$

$$\begin{cases} \text{there exist constants } \rho > 0 \text{ and } r > 0 \text{ such that } \Phi(u) \geq \rho \\ \text{for every } u \in E \text{ with } \|u\| = r \end{cases} \quad (31)$$

$$\Phi(v) \leq 0 \text{ for some } v \in E \text{ with } \|v\| > r. \quad (32)$$

Set

$$c = \inf_{P \in \mathfrak{F}} \sup_{p \in P} \Phi(p) \quad (33)$$

where  $\mathfrak{F}$  denotes the class of all paths joining 0 to  $v$ . Then there exists a sequence  $(u_j)$  in  $E$  such that  $\Phi(u_j) \rightarrow c$  and  $\Phi'(u_j) \rightarrow 0$  in  $E^*$ .

The proof of Lemma 4 is essentially the same as the proof given in [1]. In order to prove Theorems 3, 4 and 5 we apply Lemma 4 in  $E = H_0^1$  to the functional

$$\Phi(u) = \frac{1}{2} \int |\nabla u|^2 - \frac{1}{p+1} \int (u^+)^{p+1} - \int F(u^+) \quad (34)$$

where  $F(u) = \int_0^u f(t) dt$ . Property (31) is an easy consequence of assumption (24). For every  $u \in H_0^1$ ,  $u \geq 0$  in  $\Omega$ ,  $u \not\equiv 0$  we have  $\lim_{t \rightarrow +\infty} \Phi(tu) = -\infty$ . Hence, there are many  $v$ 's satisfying (32). However it is essential to make a *special choice* of  $v$  in order to be able to use properly the sequence  $(u_j)$  given by Lemma 4. More precisely we have

Lemma 5 : Under the assumptions of Theorems 3, 4 and 5 there is some  $u_0 \in H_0^1$ ,  $u_0 \geq 0$  in  $\Omega$ ,  $u_0 \not\equiv 0$  and

$$\sup_{t \geq 0} \Phi(t u_0) < \frac{1}{N} S^{N/2}. \quad (35)$$

The proof is rather technical. Let  $u_\varepsilon(x) = \frac{\phi(x)}{(\varepsilon + |x|^2)^{\frac{N-2}{2}}}$  with  $\phi \in \mathcal{D}_+(\Omega)$

and  $\phi \equiv 1$  near  $x = 0$  (assuming  $0 \in \Omega$ ). We show by an expansion method (as in Lemma 1) that  $u_\varepsilon$  satisfies (35) provided  $\varepsilon > 0$  is sufficiently small. As was already observed, the expansion technique is sensitive to the dimension  $N$ ; the cases  $N = 3$ ,  $N = 4$  and  $N \geq 5$  must be considered separately. (See the details in [5]).

Proofs of Theorem 3, 4 and 5 : By Lemma 5 there is some  $v \in H_0^1$  such that  $\|v\| > r$ ,  $\phi(v) \leq 0$  and

$$\sup_{t \geq 0} \phi(tv) < \frac{1}{N} S^{N/2} \quad (36)$$

We apply Lemma 4 with such a  $v$ . From (33) it follows that

$$c < \frac{1}{N} S^{N/2} \quad (37)$$

Let  $(u_j)$  be the sequence given by Lemma 4. We have

$$\frac{1}{2} \int |\nabla u_j|^2 - \frac{1}{p+1} \int (u_j^+)^{p+1} - \int F(u_j^+) = c + o(1) \quad (38)$$

$$-\Delta u_j = (u_j^+)^p + f(u_j^+) + \zeta_j \quad (39)$$

with  $\zeta_j \rightarrow 0$  in  $H^{-1}$ .

Combining (38) and (39) it is easy to show that  $\|u_j\|_{H^1}$  remains bounded.

Thus we may assume that

$$u_j \rightharpoonup u \quad \text{weakly in } H_0^1$$

$$u_j \rightarrow u \quad \text{a.e.}$$

From (39) we deduce that

$$-\Delta u = (u^+)^p + f(u^+) \quad \text{in } H^{-1} \quad (40)$$

By the maximum principle we have  $u \geq 0$  and so

$$-\Delta u = u^{p+1} + f(u).$$

It remains to prove that  $u \not\equiv 0$ . Suppose by contradiction that  $u \equiv 0$ . Using (23) we obtain

$$\int F(u_j^+) \rightarrow 0, \quad \int f(u_j^+) u_j \rightarrow 0 \quad (41)$$

We may always assume that

$$\int |\nabla u_j|^2 \rightarrow \ell \quad (42)$$

and by (39)

$$\int (u_j^+)^{p+1} \rightarrow \ell \quad (43)$$



From (38) we deduce that

$$\frac{1}{N} \ell = c > 0. \quad (44)$$

On the other hand we have (by Sobolev inequality)

$$\int |\nabla u_j|^2 \geq S \|u_j\|_{L^{p+1}}^2 \geq S \|u_j^+\|_{L^{p+1}}^2$$

and therefore at the limit

$$\ell \geq S \ell^{\frac{2}{p+1}}$$

Thus

$$\ell \geq S^{N/2}$$

and (by (44))

$$c \geq \frac{1}{N} S^{N/2}$$

a contradiction with (37). Therefore  $u \neq 0$ .

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