

**SUBLINEAR ELLIPTIC EQUATIONS IN  $\mathbb{R}^n$**

By

**Haïm Brezis**

and

**Shoshana Kamin**

**IMA Preprint Series # 855**

August 1991

# SUBLINEAR ELLIPTIC EQUATIONS IN $\mathbb{R}^n$

Haïm Brezis and Shoshana Kamin

## 1. Introduction

We are concerned with the question of existence (or nonexistence) and uniqueness of solutions of the problem

$$(1) \quad -\Delta u = \rho(x)u^\alpha \quad \text{in} \quad \mathbb{R}^n, \quad n \geq 3$$

with  $0 < \alpha < 1$  and  $\rho(x) \geq 0$ ,  $\rho$  not identically zero. We shall assume throughout the paper that  $\rho \in L_{loc}^\infty$ . We look for a solution  $u \geq 0$ ,  $u$  not identically zero, so that, by the strong maximum principle, if such a solution exists then  $u > 0$  in  $\mathbb{R}^n$ .

We shall often use the following:

Definition: We say that a function  $\rho \in L_{loc}^\infty(\mathbb{R}^n)$ ,  $\rho \geq 0$ , has the property (H) if the linear problem

$$(2) \quad -\Delta U = \rho \quad \text{in} \quad \mathbb{R}^n$$

has a bounded solution.

Our main result is

Theorem 1. Problem (1) has a bounded solution iff  $\rho$  satisfies (H). Moreover there is a minimal positive solution of (1).

This minimal positive solution of (1) tends to zero at infinity in a sense to be

precised later. Moreover it is the unique positive solution of (1) which tends to zero at infinity (see Theorem 2 below).

In Section 2 we prove Theorem 1 and in Section 3 we present uniqueness results for (1). In Appendix I we summarize some properties of the linear Poisson equation (2). In Appendix II we review the uniqueness question for equation (1) in bounded domains.

Problem (1) for bounded domains with zero Dirichlet condition has been extensively studied (even for more general sublinear functions). We refer in particular to Krasnoselskii [9] (Theorem 7.14 and 7.15) and [1] (see also the references therein). Problem (1) in all of space has been considered in [3], [4], and [10] under more restrictive conditions on  $\rho$  ( $\rho$  is equivalent to a radial function for large  $|x|$ ).

The study of (1) is also related to the asymptotic behavior (as  $t \rightarrow \infty$ ) of the solution of

$$(3) \quad \rho(x) \frac{\partial u}{\partial t} = \Delta u^m \quad \text{in} \quad \mathbb{R}^n \times (0, \infty)$$

with  $m = 1/\alpha > 1$  which has recently been studied by Eidus [5] (see also [6]) for a class of functions  $\rho$  tending to zero at infinity. In fact, separating variables, we have a solution  $u(x, t)$  of (3) of the form  $u(x, t) = C v^{1/m}(x)(t + \tau)^{-1/(m-1)}$  provided  $v(x)$  is a solution of (1).

### Proof of Theorem 1

#### A. Sufficient condition:

Let

$$B_R = \{x \in \mathbb{R}^n; |x| < R\}$$

and let  $u_R$  be the solution of

$$(4) \quad \begin{cases} -\Delta u = \rho u^\alpha & \text{in } B_R, \\ u = 0 & \text{on } \partial B_R. \end{cases}$$

It is well-known that  $u_R$  exists and is unique (see e.g. [9], [1] or Appendix II).

The sequence  $u_R$  is increasing with  $R$ . Indeed, let  $R' > R$ . Then  $u_{R'}$  is a supersolution for the  $R$ -problem. We now construct a subsolution  $\underline{u}$  for the  $R$ -problem with  $\underline{u} \leq u_{R'}$ . This will imply that there is a solution  $u$  for the  $R$ -problem between  $\underline{u}$  and  $u_{R'}$ . Since the unique solution is  $u_R$  it follows that  $u_R \leq u_{R'}$  in  $B_R$ . For  $\underline{u}$  we may take  $\varepsilon \varphi_1$  where  $\varphi_1$  satisfies

$$\begin{cases} -\Delta \varphi_1 = \lambda_1 \rho \varphi_1 & \text{in } B_R, \\ \varphi_1 = 0 & \text{on } \partial B_R. \end{cases}$$

We now prove that the sequence  $u_R$  remains bounded as  $R \rightarrow \infty$ . In fact

$$u_R \leq C U$$

for some appropriate constant  $C$ . Indeed,  $C U$  is a supersolution for the  $R$ -problem since

$$-\Delta (CU) = C\rho \geq \rho (CU)^\alpha$$

provided

$$C^{1-\alpha} \geq \|U\|_\infty^\alpha.$$

Therefore  $u = \lim_{R \rightarrow \infty} u_R$  exists and  $u$  is a solution of (1) satisfying

$$(5) \quad u \leq C U.$$

Clearly  $u$  is the minimal solution; indeed if  $\bar{u}$  is another solution of (1) then  $u_R \leq \bar{u}$  on  $B_R$  by the above argument and thus  $u \leq \bar{u}$ .

B. Necessary condition

Suppose  $u$  is bounded positive solution of (1) and set

$$v = \frac{1}{1 - \alpha} u^{1-\alpha}.$$

Then

$$-\Delta v = \alpha u^{-\alpha-1} |\nabla u|^2 + \rho \geq \rho.$$

The solution  $w_R$  of the problem

$$(6) \quad \begin{cases} -\Delta w_R = \rho & \text{in } B_R, \\ w_R = 0 & \text{on } \partial B_R \end{cases}$$

satisfies  $w_R \leq v$ . Thus  $w_R$  increases as  $R \rightarrow \infty$  to a bounded solution of (2).

The meaning of Theorem 1 is that if  $\rho(x)$  decays fast enough at infinity then Problem (1) has a solution. It need not exist if  $\rho(x)$  has a slow decay at infinity. As we see in the next example, if  $\rho(x)$  decays like a power, the critical exponent is two.

Example 1: Assume

$$\rho(x) = \frac{1}{1 + |x|^p} \quad \text{with } p > 2$$

or

$$\rho(x) = \frac{1}{(1 + |x|^2)|\log(2+|x|)|^p} \quad \text{with } p > 2$$

then Problem (1) has a bounded solution. Indeed the Poisson integral

$\frac{c}{|x|^{n-2}} * \rho$  provides a bounded positive solution of (2) where  $c/|x|^{n-2}$  is the fundamental solution of  $-\Delta$ .

Example 2: Assume

$$\rho(x) = \frac{1}{1 + |x|^p} \quad \text{with } p \leq 2$$

then Problem (1) has no solution. In fact a stronger nonexistence result holds.

Assume

$$(7) \quad \int_{|x| \geq 1} \frac{\rho(x)}{|x|^{n-2}} dx = \infty,$$

then there is no function  $u \in L^1_{loc}(\mathbb{R}^n)$  satisfying

$$(8) \quad \begin{cases} -\Delta u = \rho u^\alpha & \text{in } \mathcal{D}'(\mathbb{R}^n) \\ u \geq 0 \end{cases}$$

except  $u \equiv 0$ . Indeed, assume we have a solution of (8). By local regularity,  $u \in W^{2,q}_{loc}$  for all  $q < \infty$  and if  $u$  is not identically zero then  $u > 0$  in  $\mathbb{R}^n$ .

As above, set

$$v = \frac{1}{1-\alpha} u^{1-\alpha}$$

so that  $-\Delta v \geq \rho$ . It follows that

$$(9) \quad w_R \leq v$$

where  $w_R$  is defined by (6). As  $R \uparrow \infty$ ,  $w_R \uparrow \infty$  because of (7) (see Appendix I). This is impossible by (9).

**Remark 1.** The minimal solution  $u$  obtained in Theorem 1 satisfies

$$(10) \quad u(x) = c \int_{\mathbb{R}^n} \frac{\rho(y) u^\alpha(y)}{|x-y|^{n-2}} dy$$

and also

$$(11) \quad \lim_{R \rightarrow \infty} \oint_{S_R} u = 0$$

where  $\oint_{S_R} u$  denotes the average of  $u$  on the sphere of radius  $R$  (centered at 0). Indeed,  $u$  satisfies (5) for any positive solution  $U$  of (2); in particular we can take  $U = \frac{c}{|x|^{n-2}} * \rho$ . We now apply Lemma A.4 in Appendix I to conclude

that (11) holds. As a consequence of (11) we have

$$\liminf_{|x| \rightarrow \infty} u(x) = 0 .$$

Next, let  $f = \rho u^\alpha$ . The linear equation  $-\Delta v = f$  in  $\mathbb{R}^n$  has a unique solution satisfying

$$\lim_{R \rightarrow \infty} \oint_{S_R} v = 0 ,$$

namely  $v = \frac{c}{|x|^{n-2}} * f$ . Since  $u$  satisfies the same equation and also (11) we obtain (10).

**Remark 2.** The minimal solution  $u$  of (1) depends monotonically on  $\rho$ . Indeed let  $\rho_1 \leq \rho_2$  and let  $u_1, u_2$  be the corresponding minimal solutions of (1). Then  $u_2$  is a supersolution for the equation

$$\begin{aligned} -\Delta u &= \rho_1 u^\alpha & \text{in } B_R \\ u &= 0 & \text{on } \partial B_R . \end{aligned}$$

Thus  $u_{1,R} \leq u_2$  in  $B_R$ . Passing to the limit as  $R \rightarrow \infty$  we find that  $u_1 \leq u_2$ .

**Remark 3.** The minimal solution  $u$  obtained in Theorem satisfies  $C_1 U^{1-\alpha} \leq u \leq C_2 U$ . In general these bounds are sharp. For example if  $\rho$  has compact support then both  $u$  and  $U$  behave at infinity like the fundamental solution. However if  $\rho(x) \sim |x|^{-p}$  at infinity with  $2+(n-2)(1-\alpha) < p < n$  then a simple computation shows that  $U(x) \sim |x|^{-(p-2)}$  and  $u(x) \sim |x|^{-(p-2)(1-\alpha)}$ .

### 3. Uniqueness

As we have noted the minimal solution  $u$  constructed above satisfies

$$(12) \quad \liminf_{|x| \rightarrow \infty} u(x) = 0 .$$

Our main uniqueness result is

Theorem 2. Assuming  $\rho$  has property (H), then there is exactly one bounded positive solution of (1) satisfying (12).

Remark 4. There exist other bounded positive solutions of (1) which do not satisfy (12). In fact, given any positive constant  $a$ , there exists a solution of (1) satisfying

$$\liminf_{|x| \rightarrow \infty} u(x) = a.$$

Indeed, consider the problem

$$(13) \quad \begin{cases} -\Delta u = \rho u^\alpha & \text{in } B_R , \\ u = a & \text{on } \partial B_R . \end{cases}$$

As subsolution for (13) we may take  $a$  and as supersolution we may take  $(CU + a)$  where  $U = \frac{c}{|x|^{n-2}} * \rho$  with  $C$  is large enough. We then let  $R \rightarrow \infty$ .

The proof of Theorem 2 is divided into 3 steps:

Step 1. Assume  $\rho_1 \leq \rho_2$  and that they satisfy property (H). Given any bounded positive solution  $u_1$  of

$$(14) \quad \begin{cases} -\Delta u_1 = \rho_1 u_1^\alpha & \text{in } \mathbb{R}^n \\ \lim_{R \rightarrow \infty} \int_{S_R} u_1 = 0 \end{cases}$$

then there exists a bounded positive solution  $u_2$  of



$$15) \quad \begin{cases} -\Delta u_2 = \rho_2 u_2^\alpha & \text{in } \mathbb{R}^n \\ \lim_{R \rightarrow \infty} \int_{S_R} u_2 = 0 \end{cases}$$

such that  $u_1 \leq u_2$ .

Proof. Clearly  $u_1$  is a subsolution for (15) in the sense that

$$-\Delta u_1 \leq \rho_2 u_1^\alpha.$$

Since  $u_1$  is bounded we have

$$-\Delta u_1 \leq C \rho_2$$

and by Lemma A.6 we find that

$$u_1 \leq C \left( \frac{1}{|x|^{n-2}} * \rho_2 \right).$$

The right-hand side is a supersolution for (15) provided  $C$  is large enough. Using the standard monotone iteration technique (directly in  $\mathbb{R}^n$ ) we obtain a solution  $u_2$  of (15) such that

$$u_1 \leq u_2 \leq C \left( \frac{1}{|x|^{n-2}} * \rho_2 \right).$$

The only difference with the usual case of bounded domains is that the Dirichlet condition is replaced by the condition at infinity  $\lim_{R \rightarrow \infty} \int_{S_R} u = 0$ . The standard maximum principle is replaced at each stage by Lemma A.6.

We shall now show that it suffices to prove Theorem 2 in the case  $\rho > 0$ .

Step 2. Assume we have proved uniqueness for any  $\rho > 0$ , then we also have uniqueness for a general  $\rho \geq 0$ .

Proof. Let  $\rho_\varepsilon = \rho + \varepsilon h$  where  $h \in C^\infty(\mathbb{R}^n) \cap L^1(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$  with  $h > 0$ .

Let  $u_\epsilon$  be the unique solution of

$$(16) \quad \begin{cases} -\Delta u_\epsilon = \rho_\epsilon u_\epsilon^\alpha & \text{in } \mathbb{R}^n \\ \lim_{R \rightarrow \infty} \oint_{S_R} u_\epsilon = 0 \end{cases}$$

Let  $u$  be any solution of

$$(17) \quad \begin{cases} -\Delta u = \rho u^\alpha & \text{in } \mathbb{R}^n \\ \lim_{R \rightarrow \infty} \oint_{S_R} u = 0 \end{cases}.$$

By Step 1 (and by the uniqueness of  $u_\epsilon$ ) we know that

$$(18) \quad u \leq u_\epsilon$$

We prove that, as  $\epsilon \downarrow 0$ ,  $u_\epsilon \downarrow \underline{u}$  where  $\underline{u}$  is the minimal solution constructed in Theorem 1. Indeed let  $u_{\epsilon,R}$  and  $u_R$  be the positive solutions of

$$(19) \quad \begin{aligned} -\Delta u_{\epsilon,R} &= \rho_\epsilon u_{\epsilon,R}^\alpha & \text{in } B_R \\ u_{\epsilon,R} &= 0 & \text{on } \partial B_R \end{aligned}$$

and

$$(20) \quad \begin{aligned} -\Delta u_R &= \rho u_R^\alpha & \text{in } B_R \\ u_R &= 0 & \text{on } \partial B_R. \end{aligned}$$

We now use the same device as in Appendix II (method II), namely, we multiply (19) by  $u_R$  and (20) by  $u_{\epsilon,R}$ . Integrating by parts we find

$$\int_{B_R} \rho u_{\epsilon,R}^\alpha u_R^\alpha (u_{\epsilon,R}^{1-\alpha} - u_R^{1-\alpha}) = \int_{B_R} (\rho_\epsilon - \rho) u_R u_{\epsilon,R}^\alpha$$

and thus

$$\int_{B_R} \rho u_{\epsilon,R}^\alpha u_R^\alpha (u_{\epsilon,R}^{1-\alpha} - u_R^{1-\alpha}) \leq C\epsilon$$

where  $C$  is independent of  $R$ . Passing to the limit as  $R \rightarrow \infty$  (and using Fatou) we obtain

$$\int_{\mathbb{R}^n} \rho u_\varepsilon^\alpha \underline{u}^\alpha (u_\varepsilon^{1-\alpha} - \underline{u}^{1-\alpha}) \leq C\varepsilon.$$

Using (18) we have

$$\int \rho u^\alpha \underline{u}^\alpha (u^{1-\alpha} - \underline{u}^{1-\alpha}) = 0$$

and thus  $\rho u^\alpha = \rho \underline{u}^\alpha$ . Hence  $\Delta(u - \underline{u}) = 0$  and therefore  $u = \underline{u}$  (by the condition at infinity).

The last step involves the use of parabolic equations as in [8]. As we already mentioned in the Introduction if  $u(x)$  is a solution of (1) then

$$v(x, t) = \frac{Cu^{1/m}(x)}{(t + \tau)^{1/(m-1)}}$$

satisfies

$$(21) \quad \rho \frac{\partial v}{\partial t} = \Delta v^m$$

where  $m = 1/\alpha$  and  $C = (m-1)^{-1/(m-1)}$ . Our proof of uniqueness for problem (1), (12) relies heavily on existence, uniqueness and comparison properties of solution of (21).

**Step 3.** We recall first a well-known fact about bounded domains (see e.g. [2]).

Let  $\Omega \subset \mathbb{R}^n$  be a smooth bounded domain,  $\rho \in L^\infty(\Omega)$ ,  $\rho \geq \delta > 0$  on  $\Omega$ .

Then given any  $v_0 \geq 0$  on  $\Omega$ ,  $v_0 \in L^\infty(\Omega)$ , there exists a unique solution  $v(x, t)$  of the problem

$$(22) \quad \begin{cases} \rho \frac{\partial v}{\partial t} - \Delta v^m = 0 & \text{in } \Omega \times (0, \infty) \\ v = 0 & \text{on } \partial\Omega \times (0, \infty) \\ v(x, 0) = v_0(x) & \text{in } \Omega \end{cases}$$

Moreover if there is another solution  $\tilde{v}(x, t)$  of (22) with  $\tilde{v}(x, t) \geq 0$  on  $\partial\Omega \times (0, \infty)$  and  $\tilde{v}(x, 0) \geq v_0(x)$  then  $\tilde{v}(x, t) \geq v(x, t)$ .

Let  $\underline{u}$  be the minimal positive solution of (1) in the sense of Theorem 1. Let  $u$  be any bounded positive solution of (1) satisfying (12). By Appendix I we know that

$$\lim_{R \rightarrow \infty} \int_{S_R} u = 0.$$

Let  $v_R$  be the solution of

$$\begin{cases} \rho \frac{\partial v_R}{\partial t} - \Delta v_R^m = 0 & \text{in } B_R \times (0, \infty) \\ v_R(x, t) = 0 & \text{on } \partial B_R \times (0, \infty) \\ v_R(x, 0) = u(x) & \text{in } B_R \end{cases}.$$

By comparison in bounded domains we see that

$$(23) \quad v_R(x, t) \leq \frac{u^{1/m}(x)}{(t+1)^{1/(m-1)}}$$

and also

$$(24) \quad v_R(x, t) \leq \frac{u^{1/m}(x)}{t^{1/(m-1)}}.$$

As  $R \uparrow \infty$  the sequence  $v_R$  increases to some limit  $v_\infty(x, t)$  which satisfies

$$(25) \quad \rho \frac{\partial v_\infty}{\partial t} - \Delta v_\infty^m = 0 \quad \text{in } \mathbb{R}^n \times (0, \infty)$$

and

$$(26) \quad v_\infty(x, 0) = u(x).$$

Moreover we have

$$(27) \quad v_{\omega}(x, t) \leq \frac{u^{1/m}(x)}{(t+1)^{1/(m-1)}}.$$

We already have a solution of (25), (26) namely  $\frac{u^{1/m}(x)}{(t+1)^{1/(m-1)}}$ . We claim that

$$(28) \quad v_{\omega}(x, t) = \frac{u^{1/m}(x)}{(t+1)^{1/(m-1)}} \equiv \hat{v}(x, t).$$

For this purpose we multiply

$$\rho \frac{\partial}{\partial t}(\hat{v} - v_{\omega}) - \Delta(\hat{v}^m - v_{\omega}^m) = 0$$

by the function  $K(x) = c \left[ \frac{1}{|x|^{n-2}} - \frac{1}{R^{n-2}} \right]$  and integrate over  $B_R \times (0, T)$ .

We find

$$\begin{aligned} & \int_{B_R} \rho(x) (\hat{v} - v_{\omega}) K(x) dx \Big|_{t=T} + \int_0^T (\hat{v}^m - v_{\omega}^m) dt \Big|_{x=0} \\ &= - \int_0^T \int_{\partial B_R} (\hat{v}^m - v_{\omega}^m) \frac{\partial K}{\partial \nu} dS dt. \end{aligned}$$

The integral on the right hand side is bounded by

$$C T \int_{S_R} u$$

which tends to zero as  $R \rightarrow \infty$ . Thus  $\hat{v} = v_{\omega}$  (since  $\rho > 0$ ). Passing to the limit in (24) we find

$$\frac{u^{1/m}(x)}{(t+1)^{1/(m-1)}} \leq \frac{u^{1/m}(x)}{t^{1/(m-1)}}.$$

Letting  $t \rightarrow \infty$  we conclude that  $u \leq \underline{u}$ .

Remark 5. Assume  $\rho$  has property (H). As we know from Appendix I

$$\lim_{R \rightarrow \infty} \oint_{S_R} U = 0$$

where  $U = \frac{c}{|x|^{n-2}} * \rho$ , and thus  $\lim_{|x| \rightarrow \infty} \inf U = 0$ . It may happen that  $U(x)$  does not tend to zero as  $|x| \rightarrow \infty$ . Here is a simple example for  $n \geq 4$ . Let  $\psi(x')$  be the solution of

$$\begin{cases} -\Delta_{x'} \psi = \rho(x') & \text{in } \mathbb{R}^{n-1} \\ \lim_{|x'| \rightarrow \infty} \psi = 0 \end{cases}$$

where  $\rho \in C_0^\infty(\mathbb{R}^{n-1})$ ,  $\rho \geq 0$  and  $\rho$  not identically zero. Then

$$U(x) = \psi(x') \quad x = (x_1, x')$$

provides such an example since  $U(x_1, 0) = \psi(0)$  does not tend to zero as  $|x_1| \rightarrow \infty$ . In such a situation there is no solution  $u$  of (1) which tends to zero at infinity because of the estimate from below  $u^{1-\alpha} \geq (1-\alpha)U$  (see the proof of necessary condition in Theorem 1).

The uniqueness question becomes easier under a stronger assumption

Theorem 2'. Assume there is a solution  $U$  of (2) such that

$$(29) \quad \lim_{|x| \rightarrow \infty} U(x) = 0.$$

Then there exists a unique positive solution  $u$  of (1) such that

$$\lim_{|x| \rightarrow \infty} u(x) = 0.$$

Proof. The existence part is clear since we already know that there is a solution  $u$  of (1) such that  $u \leq CU$ . For the uniqueness we could invoke Theorem 2 but we present instead a simple argument due to Louis Nirenberg.

First we change the unknown. As in the proof of Theorem 1 we set

$$v = \frac{1}{1-\alpha} u^{1-\alpha}$$

so that we find

$$(30) \quad -\Delta v - \frac{C}{v} |\nabla v|^2 = \rho$$

for some positive constant  $C$  (depending on  $\alpha$ ). Uniqueness holds for (30) since the function  $1/v$  is decreasing in  $v$ . More precisely, suppose we have two solutions  $v_1, v_2$  of (30) with  $\lim_{|x| \rightarrow \infty} v_1 = \lim_{|x| \rightarrow \infty} v_2 = 0$ . Then  $w = v_1 - v_2$  satisfies

$$-\Delta w - \frac{C}{v_1} \nabla(v_1 + v_2) \cdot \nabla w + \frac{C}{v_1 v_2} |\nabla v_2|^2 w = 0.$$

Since the coefficient of  $w$  is nonnegative we may use the maximum principle to conclude that  $w = 0$ .

**Remark 6.** Clearly if  $\rho$  is a radial function satisfying (H) then (29) holds. It also holds if  $\rho$  is bounded by a radial function satisfying (H).

#### 4. Some generalization

Our methods extend to more general problems of the form

$$-\Delta u = \rho(x) f(u) \quad \text{in } \mathbb{R}^n$$

under suitable assumptions of  $f$  and in particular  $f(u)$  behaves like  $u^\alpha$  near  $u = 0$ . For simplicity we restrict our attention to the model problem  $f(u) = u^{\alpha(1-u)}$ , i.e.

$$(31) \quad -\Delta u = \rho(x) u^{\alpha(1-u)} \quad \text{in } \mathbb{R}^n.$$

**Theorem 3.** Assume  $\rho$  satisfies (H). Then there is a unique solution  $u$ ,  $0 < u < 1$  of (31) such that

$$(32) \quad \lim_{|x| \rightarrow \infty} \inf u(x) = 0 .$$

Proof. For the existence part we proceed as in the proof of Theorem 1 (sufficient condition). We obtain a minimal solution  $\underline{u}$  with  $\underline{u} \leq 1$  and  $\underline{u} \leq CU$ . For the uniqueness we proceed in two steps.

Step 1. Let  $u$  be any solution of (31), (32). Then there exists some  $\varepsilon > 0$  such that

$$(33) \quad \varepsilon u \leq \underline{u} .$$

It is useful to introduce the unique positive solution  $v$  of the problem

$$(34) \quad \begin{cases} -\Delta v = \rho v^\alpha & \text{in } \mathbb{R}^n \\ \oint_{S_R} v \rightarrow 0 & \text{as } R \rightarrow \infty \end{cases} .$$

Note that  $u$  is a subsolution for (32) since

$$u^{\alpha(1-u)} \leq u^\alpha$$

and therefore, by monotone iteration and uniqueness of  $v$ , we obtain

$$u \leq v .$$

Next, we note that for  $\varepsilon > 0$  small enough  $\varepsilon v$  is a subsolution for (31) since

$$-\Delta(\varepsilon v) = \varepsilon \rho v^\alpha \leq \rho(\varepsilon v)^\alpha (1 - \varepsilon v) .$$

It follows that  $\varepsilon v \leq \underline{u}$ , the minimal solution of (31) (to justify this we use comparison in  $B_R$  and then let  $R \rightarrow \infty$ ). Thus (33) holds.

Step 2. We now follow the same technique as in Method III of Appendix II. Let  $u$  be any solution of (31), (32) and let

$$\Lambda = \{t \in [0, 1]; \quad tu \leq \underline{u}\} .$$



We claim that  $1 \in \Lambda$ . Suppose not, that

$$t_0 = \sup \Lambda < 1.$$

By Step 1 we know that  $t_0 > 0$ . Fix  $K$  large enough so that the function  $f(t) + Kt$  is increasing on  $[0, 1]$ . We have

$$-\Delta(\underline{u} - t_0 u) + K\rho(\underline{u} - t_0 u) \geq \rho[f(t_0 u) - t_0 f(u)].$$

Note that for  $\varepsilon > 0$  small enough

$$f(t_0 u) - t_0 f(u) \geq \varepsilon f(u)$$

since

$$t_0^\alpha \geq t_0 + \varepsilon.$$

Thus we obtain

$$-\Delta(\underline{u} - t_0 u - \varepsilon u) \geq 0$$

and by Appendix I we conclude that  $\underline{u} - t_0 u - \varepsilon u \geq 0$ . Hence  $t_0 + \varepsilon \in \Lambda$ , which contradicts the maximality of  $t_0$ .

### Appendix I

Throughout the paper we have often used the property (H), namely that the equation

$$(A.1) \quad -\Delta U = f \quad \text{in } \mathbb{R}^n$$

has a bounded solution. We discuss here some equivalent forms and some consequences. In what follows we always assume that  $f \in L_{loc}^\infty(\mathbb{R}^n)$ ,  $f \geq 0$  a.e. and that  $f$  is not identically zero. Let  $u_R$  be the solution of

$$(A.2) \quad \begin{cases} -\Delta u_R = f & \text{in } B_R \\ u_R = 0 & \text{on } \partial B_R \end{cases}.$$

Note that  $u_R$  is a nondecreasing sequence of positive functions (in  $B_R$ ) for  $R$  large enough. Moreover  $u_R$  is given by

$$(A.3) \quad u_R(x) = \int_{B_R} G_R(x,y) f(y) dy$$

where  $G_R$  is the Green's function relative to  $B_R$  and zero boundary condition. Let

$$u_\infty(x) = \lim_{R \uparrow \infty} u_R(x) \quad (\text{possibly } +\infty).$$

Note that, by monotone convergence of  $G_R$ ,

$$u_\infty(x) = c \int_{\mathbb{R}^n} \frac{1}{|x-y|^{n-2}} f(y) dy = \frac{c}{|x|^{n-2}} * f$$

(possibly  $+\infty$ ), where  $c/|x|^{n-2}$  is the fundamental solution. Remark that there are only two possibilities, either  $u_\infty(x) = +\infty \quad \forall x$  or  $u_\infty(x) < +\infty \quad \forall x$ . Indeed suppose for example that  $u_\infty(0) < +\infty$ .

Write

$$u_\infty(x) = c \int_{|y| \leq 2|x|} \frac{f(y)}{|x-y|^{n-2}} + c \int_{|y| > 2|x|} \frac{f(y)}{|x-y|^{n-2}}.$$

The first integral is finite (for each fixed  $x$ ) while the second integral is bounded by  $2^{n-2} c \int \frac{f(y)}{|y|^{n-2}} dy$ . Hence  $u_\infty(x) < \infty$ .

If we make the assumption that

$$u_\infty(0) = c \int \frac{f(y)}{|y|^{n-2}} dy < \infty$$

then  $u_\infty(x)$  is finite for each fixed  $x$  but it need not be uniformly bounded on  $\mathbb{R}^n$ .

Lemma A.1.  $f$  satisfies property (H) iff

$$(A.4) \quad \frac{c}{|x|^{n-2}} * f \in L^\infty(\mathbb{R}^n) .$$

Proof. Suppose first that (H) holds. By adding a constant we may always assume that  $U \geq 0$  in  $\mathbb{R}^n$ . By the maximum principle

$$u_R \leq U \quad \text{on } B_R$$

and therefore

$$(A.5) \quad u_\infty = \frac{c}{|x|^{n-2}} * f \leq U .$$

Conversely, the function  $\frac{c}{|x|^{n-2}} * f$  provides a bounded solution of (A.1).

Since  $U$  could be any nonnegative solution of (A.1) we have

Corollary A.2.  $u_\infty$  is the minimal positive solution of (A.1).

As a consequence of minimality we have

Corollary A.3.

$$\lim_{|x| \rightarrow \infty} \inf u_\infty(x) = 0 .$$

In fact, any bounded solution  $U$  of (A.1) such that

$$\lim_{|x| \rightarrow \infty} \inf U(x) = 0$$

coincides with  $u_\infty$ . This follows from the fact that the difference of any two bounded solutions of (A.1) is a bounded harmonic function and thus it is a constant.

A stronger way of expressing that  $u_\infty$  tends to zero at infinity is the following

Lemma A.4.

$$\lim_{R \rightarrow \infty} \oint_{S_R} u_{\infty} = 0$$

where  $\oint_{S_R}$  denotes the average on the sphere of radius  $R$ .

Proof. By Fubini we have

$$\frac{1}{R^{n-1}} \int_{S_R} u_{\infty}(y) dS_y = c \int_{\mathbb{R}^n} f(x) \frac{1}{R^{n-1}} \left[ \int_{|y|=R} \frac{dS_y}{|x-y|^{n-2}} \right] dx .$$

Note that

$$I(x) = \int_{|x|=R} \frac{dS_y}{|x-y|^{n-2}} = \begin{cases} CR \left( \frac{R}{|x|} \right)^{n-2} & \text{if } |x| > R \\ I(0) & \text{if } |x| < R \end{cases}$$

with

$$I(0) = \int_{|y|=R} \frac{dS_y}{|y|^{n-2}} = CR$$

(this is a consequence of the fact that  $I(x)$  is harmonic in  $|x| < R$  and in  $|x| > R$ ; moreover  $I(x) = I(|x|)$  and in addition  $I(\infty) = 0$ ).

Hence we have

$$\oint_{S_R} u_{\infty} = \frac{C}{R^{n-2}} \int_{|x| < R} f(x) dx + c \int_{|x| > R} \frac{f(x)}{|x|^{n-2}} dx .$$

Clearly the second integral tends to zero as  $R \rightarrow \infty$ . We estimate the first one by

$$\frac{C}{R^{n-2}} \int_{|x| < R_0} f(x) dx + C \int_{R_0 < |x| < R} \frac{f(x)}{|x|^{n-2}} dx .$$

We first choose  $R_0$  so that

$$C \int_{R_0 < |x|} \frac{f(x)}{|x|^{n-2}} dx < \varepsilon$$

and then  $R$  large enough so that

$$\frac{C}{R^{n-2}} \int_{|x| < R_0} f(x) dx < \varepsilon .$$

Lemma A.5. Any bounded solution  $U$  of (A.1) such that

$$\oint_{S_R} U \longrightarrow 0 \quad \text{as } R \longrightarrow \infty$$

coincides with  $u_\infty$ .

This is clear since the difference of two bounded solutions of (A.1) is a constant.

Lemma A.6. Assume  $U \in L^\infty$  with  $\Delta U \in L^\infty_{loc}$  satisfies

$$-\Delta U \leq \rho \quad \text{in } \mathbb{R}^n$$

and

$$\oint_{S_R} U \longrightarrow 0 \quad \text{as } R \longrightarrow \infty .$$

Then  $U \leq u_\infty$ .

Proof. Set  $g = -\Delta(u_\infty - U) \geq 0$ .

Since  $\oint_{S_R} (u_\infty - U) \longrightarrow 0 \quad \text{as } R \longrightarrow \infty$

we may apply Lemma A.5 to conclude that

$$u_{\infty} - U = \frac{c}{|x|^{n-2}} * g \geq 0.$$

## Appendix II

Here we briefly review several proofs of uniqueness for the problem

$$(A.6) \quad \begin{cases} -\Delta u = \rho(x) f(u) & \text{in } \Omega \\ u > 0 & \text{in } \Omega \\ u = \psi \geq 0 & \text{on } \partial\Omega \end{cases}$$

under the assumptions that  $\frac{f(t)}{t}$  is decreasing,  $\Omega$  is a smooth bounded domain and  $\rho \geq 0$ .

Method I. This is the method introduced in [1]. Let  $u_1$  and  $u_2$  be two solutions of (A.6). We have

$$(A.7) \quad -\frac{\Delta u_1}{u_1} + \frac{\Delta u_2}{u_2} = \rho \left( \frac{f(u_1)}{u_1} - \frac{f(u_2)}{u_2} \right)$$

Multiplying (A.7) by  $(u_1^2 - u_2^2)$  we obtain

$$\int \left| \nabla u_1 - \frac{u_1}{u_2} \nabla u_2 \right|^2 + \left| \nabla u_2 - \frac{u_2}{u_1} \nabla u_1 \right|^2 = \int \rho \left( \frac{f(u_1)}{u_1} - \frac{f(u_2)}{u_2} \right) (u_1^2 - u_2^2) .$$

It follows that  $u_1 = u_2$  on the set  $[\rho > 0]$ . In particular  $\rho f(u_1) = \rho f(u_2)$  on  $\Omega$ . Going back to (A.6) we see that  $u_1 = u_2$ .

Method II. Let  $u_1$  and  $u_2$  be two solution of (A.6). We have

$$(A.8) \quad -(\Delta u_1) u_2 + (\Delta u_2) u_1 = \rho u_1 u_2 \left( \frac{f(u_1)}{u_1} - \frac{f(u_2)}{u_2} \right).$$

Integrating (A.8) on the set  $[u_1 > u_2] = E$  we obtain formally

$$-\int_{\partial E} \frac{\partial u_1}{\partial \nu} u_2 + \int_{\partial E} \frac{\partial u_2}{\partial \nu} u_1 = \int_E \rho u_1 u_2 \left( \frac{f(u_1)}{u_1} - \frac{f(u_2)}{u_2} \right)$$

Note that  $u_1 = u_2$  and  $\frac{\partial}{\partial \nu}(u_1 - u_2) \leq 0$  on  $\partial E$ . Thus the lefthand side is nonnegative while the integrand on the righthand side is nonpositive. Similarly, using  $F = [u_1 < u_2]$ , we are led to

$$\int_{\Omega} \rho u_1 u_2 \left| \frac{f(u_1)}{u_1} - \frac{f(u_2)}{u_2} \right| = 0.$$

We conclude as above.

To make this argument rigorous we proceed as follows. Let  $\theta$  be a smooth nondecreasing function such that  $\theta(0) = 0$  and  $\theta(t) = 1$  for  $t \geq 1$ ,  $\theta(t) = -1$  for  $t \leq -1$ . Set

$$\theta_\varepsilon(t) = \theta(t/\varepsilon)$$

Multiplying (A.8) by  $\theta_\varepsilon(u_1 - u_2)$  and integrating we obtain

$$(A.9) \quad \begin{cases} \int [(\nabla u_1) \cdot u_2 - (\nabla u_2) \cdot u_1] \theta'_\varepsilon(u_1 - u_2) \cdot \nabla(u_1 - u_2) \\ = \int \rho u_1 u_2 \left( \frac{f(u_1)}{u_1} - \frac{f(u_2)}{u_2} \right) \theta_\varepsilon(u_1 - u_2) \end{cases}.$$

Clearly 
$$\text{LHS} \geq \int (\nabla u_2)(u_2 - u_1) \theta'_\epsilon(u_1 - u_2) \cdot \nabla(u_1 - u_2) \cdot$$

Note that

$$\int \nabla u_2(u_2 - u_1) \theta'_\epsilon(u_1 - u_2) \nabla(u_1 - u_2) = - \int \nabla u_2 \nabla \gamma_\epsilon(u_1 - u_2)$$

where 
$$\gamma_\epsilon(t) = \int_0^t s \theta'_\epsilon(s) ds .$$

Since  $|\gamma_\epsilon(t)| \leq C \epsilon$  and  $\Delta u_2 \in L^\infty$  we see that

$$\text{LHS} \geq - C \epsilon .$$

Going back to (A.9) we obtain, as  $\epsilon \rightarrow 0$ ,

$$\int \rho u_1 u_2 \left| \frac{f(u_1)}{u_1} - \frac{f(u_2)}{u_2} \right| = 0 .$$

Method III. This is a variant of Krasnoselkii's method [9]. Let  $u_1$  and  $u_2$  be two solutions. Let

$$\Lambda = \{t \in [0,1]; \quad t u_1 \leq u_2 \quad \text{on} \quad \Omega\} .$$

Clearly  $\Lambda$  contains a neighbourhood of 0. We claim that  $1 \in \Lambda$ . Suppose not, that

$$t_0 = \sup \Lambda < 1 .$$

Then

$$- \Delta(u_2 - t_0 u_1) = \rho f(u_2) - t_0 \rho f(u_1).$$

Fix a positive constant  $K$  large enough so that  $f(t) + Kt$  is increasing on  $[0, \text{Max } u_2]$ . Then

$$\begin{aligned} - \Delta(u_2 - t_0 u_1) + K \rho(u_2 - t_0 u_1) &= \rho[f(u_2) + K u_2 - t_0(f(u_1) + K u_1)] \\ &\geq \rho[f(t_0 u_1) + K t_0 u_1 - t_0(f(u_1) + K u_1)] = \rho[f(t_0 u_1) - t_0 f(u_1)] \geq 0 \end{aligned}$$



(the last inequality follows from the fact that  $f(u)/u$  is decreasing). On  $\partial\Omega$  we have  $u_2 - t_0 u_1 = (1 - t_0) \varphi \geq 0$ .

We distinguish two cases:

Case 1:  $\varphi \equiv 0$ . Using the strong maximum principle we see that either  $u_2 - t_0 u_1 > 0$  on  $\Omega$  with  $\frac{\partial}{\partial \nu}(u_2 - t_0 u_1) < 0$  on  $\partial\Omega$ . Then, clearly there is some  $\varepsilon > 0$  such that  $u_2 - t_0 u_1 \geq \varepsilon u_1$ . Thus  $t_0 + \varepsilon \in \Lambda$ . Impossible. Or  $u_2 - t_0 u_1 \equiv 0$ . This case is also impossible since we would have, by the equation  $\rho f(u_2) = t_0 \rho f(u_1)$ , but  $f(t_0 u_1) > t_0 f(u_1)$ .

Case 2:  $\varphi$  is not identically zero. We claim that there is some  $\varepsilon > 0$  such that

$$w \equiv u_2 - t_0 u_1 \geq \varepsilon u_1.$$

Suppose not, that for every  $\varepsilon > 0$  there is some point  $x_\varepsilon \in \bar{\Omega}$  such that

$$w(x_\varepsilon) < \varepsilon u_1(x_\varepsilon).$$

Clearly  $x_\varepsilon \notin \partial\Omega$  (for  $\varepsilon$  small). Choosing a point of minimum for the function  $(w - \varepsilon u_1)$  we may also assume that

$$\nabla w(x_\varepsilon) = \varepsilon \nabla u_1(x_\varepsilon).$$

As  $\varepsilon \rightarrow 0$  (through an appropriate sequence)  $x_\varepsilon \rightarrow x_0 \in \bar{\Omega}$  such that

$$w(x_0) \leq 0 \quad \text{and} \quad \nabla w(x_0) = 0.$$

It follows that  $w(x_0) = 0$  and thus  $x_0 \in \partial\Omega$ . This contradicts the strong

maximum principle since we have

$$\begin{cases} -\Delta w + K\rho w \geq 0 & \text{in } \Omega, \\ w \geq 0 & \text{on } \partial\Omega, \\ w \text{ not identically zero.} \end{cases}$$

Method IV. This is a variant of Nirenberg's method already presented in the proof of Theorem 2'. It requires further restrictions on  $f$ , namely,  $f$  is positive,

concave and  $\int_0^\delta \frac{dt}{f(t)} < \infty$ .

We use the new unknown

$$v = \int_0^u \frac{dt}{f(t)},$$

or in other words  $u = h(v)$  where  $h$  satisfies

$$h'(s) = f(h(s)).$$

The equation for  $v$  becomes

$$-\Delta v - f'(h(v))|\nabla v|^2 = \rho.$$

Uniqueness holds provided the function  $f'(h(v))$  is nonincreasing in  $v$  (see the proof of Theorem 2'). This follows from the assumptions on  $f$ .

Acknowledgements. We thank H. Berestycki, E. Gluskin, A. Edelson, H. Egnell, D. Eidus, R. Kersner and L. Nirenberg for useful discussions.

Part of this paper was written while both authors were visiting IMA at the University of Minnesota. The second author also thanks the Université Paris VI and Rutgers University for their hospitality.

### References

- [1] H. Brezis – L. Oswald, Remarks on sublinear elliptic equations, *Nonlinear Analysis*, 10 (1986), p. 55–64.
- [2] E. DiBenedetto, Continuity of weak solutions to a general porous medium equation, *Indiana University, Math. J.* 32 (1983), p. 83–118.
- [3] A.L. Edelson, Asymptotic properties of semilinear equations, *Can. Math. Bull* 32 (1989), p. 34–46.
- [4] H. Egnell, Asymptotic results for finite energy solutions of semilinear elliptic equations (to appear).
- [5] D. Eidus, The Cauchy problem for the nonlinear filtration equation in an inhomogeneous medium, *J. Diff. Eq.* 84 (1990), p. 309–318.
- [6] D. Eidus – S. Kamin, in preparation.
- [7] P. Hess, On the uniqueness of positive solutions of nonlinear elliptic boundary value problems, *Math. Z.* 154 (1977), p. 17–18.
- [8] S. Kamin – P. Rosenau, Nonlinear thermal evolution in an inhomogeneous medium, *J. Math. Phys.* 23 (1982), p. 1385–1390.
- [9] M. Krasnoselskii, Positive solutions of operator equations, Noordhoff (1964).
- [10] M. Naito, A note on bounded positive entire solutions of semilinear elliptic equations, *Hiroshima Math. J.* 14 (1984), p. 211–214.

Haïm Brezis  
 Université Paris IV  
 4, pl. Jussieu  
 75252 Paris Cedex 05  
 and  
 Rutgers University  
 New Brunswick, NJ 08903

Shoshana Kamin  
 Raymond and Beverly  
 Sackler Faculty of Exact Sciences  
 Tel-Aviv University  
 Tel-Aviv, Israel

#	Author/s	Title
774	L.A. Peletier & W.C. Troy,	Self-similar solutions for infiltration of dopant into semiconductors
775	H. Scott Dumas and James A. Ellison,	Nekhoroshev's theorem, ergodicity, and the motion of energetic charged particles in crystals
776	Stathis Filippas and Robert V. Kohn,	Refined asymptotics for the blowup of $u_t - \Delta u = u^p$ .
777	Patricia Bauman, Nicholas C. Owen and Daniel Phillips,	Maximum principles and a priori estimates for an incompressible material in nonlinear elasticity
778	Patricia Bauman, Nicholas C. Owen and Daniel Phillips,	Maximal smoothness of solutions to certain Euler–Lagrange equations from nonlinear elasticity
779	Jack Carr and Robert Pego,	Self-similarity in a coarsening model in one dimension
780	J.M. Greenberg,	The shock generation problem for a discrete gas with short range repulsive forces
781	George R. Sell and Mario Taboada,	Local dissipativity and attractors for the Kuramoto–Sivashinsky equation in thin 2D domains
782	T. Subba Rao,	Analysis of nonlinear time series (and chaos) by bispectral methods
783	Nicholas Baumann, Daniel D. Joseph, Paul Mohr and Yuriko Renardy,	Vortex rings of one fluid in another free fall
784	Oscar Bruno, Avner Friedman and Fernando Reitich,	Asymptotic behavior for a coalescence problem
785	Johannes C.C. Nitsche,	Periodic surfaces which are extremal for energy functionals containing curvature functions
786	F. Abergel and J.L. Bona,	A mathematical theory for viscous, free-surface flows over a perturbed plane
787	Gunduz Caginalp and Xinfu Chen,	Phase field equations in the singular limit of sharp interface problems
788	Robert P. Gilbert and Yongzhi Xu,	An inverse problem for harmonic acoustics in stratified oceans
789	Roger Fosdick and Eric Volkmann,	Normality and convexity of the yield surface in nonlinear plasticity
790	H.S. Brown, I.G. Kevrekidis and M.S. Jolly,	A minimal model for spatio-temporal patterns in thin film flow
791	Chao–Nien Chen,	On the uniqueness of solutions of some second order differential equations
792	Xinfu Chen and Avner Friedman,	The thermistor problem for conductivity which vanishes at large temperature
793	Xinfu Chen and Avner Friedman,	The thermistor problem with one-zero conductivity
794	E.G. Kalnins and W. Miller, Jr.,	Separation of variables for the Dirac equation in Kerr Newman space time
795	E. Knobloch, M.R.E. Proctor and N.O. Weiss,	Finite-dimensional description of doubly diffusive convection
796	V.V. Pukhnachov,	Mathematical model of natural convection under low gravity
797	M.C. Knaap,	Existence and non-existence for quasi-linear elliptic equations with the p-laplacian involving critical Sobolev exponents
798	Stathis Filippas and Wenxiong Liu,	On the blowup of multidimensional semilinear heat equations
799	A.M. Meirmanov,	The Stefan problem with surface tension in the three dimensional case with spherical symmetry: non-existence of the classical solution
800	Bo Guan and Joel Spruck,	Interior gradient estimates for solutions of prescribed curvature equations of parabolic type
801	Hi Jun Choe,	Regularity for solutions of nonlinear variational inequalities with gradient constraints
802	Peter Shi and Yongzhi Xu,	Quasistatic linear thermoelasticity on the unit disk
803	Satyanad Kichenassamy and Peter J. Olver,	Existence and non-existence of solitary wave solutions to higher order model evolution equations
804	Dening Li,	Regularity of solutions for a two-phase degenerate Stefan Problem
805	Marek Fila, Bernhard Kawohl and Howard A. Levine,	Quenching for quasilinear equations
806	Yoshikazu Giga, Shun'ichi Goto and Hitoshi Ishii,	Global existence of weak solutions for interface equations coupled with diffusion equations
807	Mark J. Friedman and Eusebius J. Doedel,	Computational methods for global analysis of homoclinic and heteroclinic orbits: a case study
808	Mark J. Friedman,	Numerical analysis and accurate computation of heteroclinic orbits in the case of center manifolds
809	Peter W. Bates and Songmu Zheng,	Inertial manifolds and inertial sets for the phase-field equations
810	J. López Gómez, V. Márquez and N. Wolanski,	Global behavior of positive solutions to a semilinear equation with a nonlinear flux condition
811	Xinfu Chen and Fahuai Yi,	Regularity of the free boundary of a continuous casting problem
812	Eden, A., Foias, C., Nicolaenko, B. and Temam, R.,	Inertial sets for dissipative evolution equations Part I: Construction and applications
813	Jose–Francisco Rodrigues and Boris Zaltzman,	On classical solutions of the two-phase steady-state Stefan problem in strips
814	Viorel Barbu and Srdjan Stojanovic,	Controlling the free boundary of elliptic variational inequalities on a variable domain
815	Viorel Barbu and Srdjan Stojanovic,	A variational approach to a free boundary problem arising in electro-photography
816	B.H. Gilding and R. Kersner,	Diffusion-convection-reaction, free boundaries, and an integral equation

- 817 Shoshana Kamin, Lambertus A. Peletier and Juan Luis Vazquez, On the Barenblatt equation of elasto-plastic filtration
- 818 Avner Friedman and Bei Hu, The Stefan problem with kinetic condition at the free boundary
- 819 M.A. Grinfeld, The stress driven instabilities in crystals: mathematical models and physical manifestations
- 820 Bei Hu and Lihe Wang, A free boundary problem arising in electrophotography: solutions with connected toner region
- 821 Yongzhi Xu, T. Craig Poling, and Trent Brundage, Direct and inverse scattering of time harmonic acoustic waves in an inhomogeneous shallow ocean
- 822 Steven J. Altschuler, Singularities of the curve shrinking flow for space curves
- 823 Steven J. Altschuler and Matthew A. Grayson, Shortening space curves and flow through singularities
- 824 Tong Li, On the Riemann problem of a combustion model
- 825 L.A. Peletier & W.C. Troy, Self-similar solutions for diffusion in semiconductors
- 826 C.J. van Duijn, L.A. Peletier & R.J. Schotting, On the analysis of brine transport in porous media
- 827 Minkyu Kwak, Finite dimensional description of convective reaction-diffusion equations
- 828 Minkyu Kwak, Finite dimensional inertial forms for the 2D Navier–Stokes equations
- 829 Victor A. Galaktionov and Sergey A. Posashkov, On some monotonicity in time properties for a quasilinear parabolic equation with source
- 830 Victor A. Galaktionov, Remark on the fast diffusion equation in a ball
- 831 Hi Jun Choe and Lihe Wang, A regularity theory for degenerate vector valued variational inequalities
- 832 Vladimir I. Olikar and Nina N. Uraltseva, Evolution of nonparametric surfaces with speed depending on curvature, II. The mean curvature case.
- 833 S. Kamin and W. Liu, Large time behavior of a nonlinear diffusion equation with a source
- 834 Shoshana Kamin and Juan Luis Vazquez, Singular solutions of some nonlinear parabolic equations
- 835 Bernhard Kawohl and Robert Kersner, On degenerate diffusion with very strong absorption
- 836 Avner Friedman and Fernandor Reitich, Parameter identification in reaction-diffusion models
- 837 E.G. Kalnins, H.L. Manocha and Willard Miller, Jr., Models of  $q$ -algebra representations I. Tensor products of special unitary and oscillator algebras
- 838 Robert J. Sacker and George R. Sell, Dichotomies for linear evolutionary equations in Banach spaces
- 839 Oscar P. Bruno and Fernando Reitich, Numerical solution of diffraction problems: a method of variation of boundaries
- 840 Oscar P. Bruno and Fernando Reitich, Solution of a boundary value problem for Helmholtz equation via variation of the boundary into the complex domain
- 841 Victor A. Galaktionov and Juan L. Vazquez, Asymptotic behaviour for an equation of superslow diffusion. The Cauchy problem
- 842 Josephus Hulshof and Juan Luis Vazquez, The Dipole solution for the porous medium equation in several space dimensions
- 843 Shoshana Kamin and Juan Luis Vazquez, The propagation of turbulent bursts
- 844 Miguel Escobedo, Juan Luis Vazquez and Enrike Zuazua, Source-type solutions and asymptotic behaviour for a diffusion-convection equation
- 845 Marco Biroli and Umberto Mosco, Discontinuous media and Dirichlet forms of diffusion type
- 846 Stathis Filippas and Jong-Shenq Guo, Quenching profiles for one-dimensional semilinear heat equations
- 847 H. Scott Dumas, A Nekhoroshev-like theory of classical particle channeling in perfect crystals
- 848 R. Natalini and A. Tesei, On a class of perturbed conservation laws
- 849 Paul K. Newton and Shinya Watanabe, The geometry of nonlinear Schrödinger standing waves
- 850 S.S. Sritharan, On the nonsmooth verification technique for the dynamic programming of viscous flow
- 851 Mario Taboada and Yuncheng You, Global attractor, inertial manifolds and stabilization of nonlinear damped beam equations
- 852 Shigeru Sakaguchi, Critical points of solutions to the obstacle problem in the plane
- 853 F. Abergel, D. Hilhorst and F. Issard-Roch, On a dissolution-growth problem with surface tension in the neighborhood of a stationary solution
- 854 Erasmus Langer, Numerical simulation of MOS transistors
- 855 Haim Brezis and Shoshana Kamin, Sublinear elliptic equations in  $\mathbb{R}^n$
- 856 Johannes C.C. Nitsche, Boundary value problems for variational integrals involving surface curvatures
- 857 Chao–Nien Chen, Multiple solutions for a semilinear elliptic equation on  $\mathbb{R}^N$  with nonlinear dependence on the gradient
- 858 D. Brochet, X. Chen and D. Hilhorst, Finite dimensional exponential attractor for the phase field model
- 859 Joseph D. Fehribach, Mullins-Sekerka stability analysis for melting-freezing waves in helium-4
- 860 Walter Schempp, Quantum holography and neurocomputer architectures
- 861 D.V. Anosov, An introduction to Hilbert’s 21st problem
- 862 Herbert E Huppert and M Grae Worster, Vigorous motions in magma chambers and lava lakes