EXISTENCE OF GROUND STATES FOR A QUASILINEAR COUPLED SYSTEM IN \mathbb{R}^N

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ABSTRACT. In this work we consider the following class of quasilinear coupled systems

 $(S_{\theta}) \left\{ \begin{array}{ll} -\Delta u + a(x)u - \Delta(u^2)u = g(u) + \theta\lambda(x)|u|^{\alpha-2}u|v|^{\beta}, & x \in \mathbb{R}^N, \\ -\Delta v + b(x)v - \Delta(v^2)v = h(v) + \theta\lambda(x)|v|^{\beta-2}v|u|^{\alpha}, & x \in \mathbb{R}^N, \end{array} \right.$

where $N \geq 3$ and $a, b: \mathbb{R}^N \to \mathbb{R}$ are positive potentials, $\lambda: \mathbb{R}^N \to \mathbb{R}$ is a suitable continuous function, $\theta > 0$ and $\alpha, \beta > 2$ satisfying $\alpha + \beta < 2.2^*$. On the nonlinear terms we assume that g, h are in C^1 class and are superlinear functions at infinity and at the origin. The main theorem is stated without the well known Ambrosetti-Rabinowitz condition at infinity. Using a change of variable, we turn the quasilinear coupled system into a nonlinear coupled system, where we establish a variational approach based on Nehari method. 1. INTRODUCTION

We look for ground states for the general class of quasilinear coupled systems involving Schrödinger equations (S_{θ}) . This class of systems imposes some difficulties. The first one is that the energy functional associated to System (S_{θ}) is not well defined in the whole space $H^1(\mathbb{R}^N)^2$. Thus, motivated by seminal works [1–6] we also use a change of variable to reformulate our initial problem, obtaining a nonlinear coupled system. After change of variable, the modified problem has an associated energy functional well defined in the whole space $H^1(\mathbb{R}^N)^2$ and the solutions are related with solutions of the initial System (S_{θ}) . The second difficulty is the lack of compactness due to the fact that the system is defined in the whole Euclidean space \mathbb{R}^N . Moreover, System (S_{θ}) involve strongly coupled Schrödinger equations because of the coupling terms in the right hand side. We emphasize that we do not use the well known Ambrosetti-Rabinowitz condition. We suppose that the potentials a, b satisfy the following hypotheses:

- (a_0) $a, b, \lambda \in C(\mathbb{R}^N, \mathbb{R})$ are 1-periodic functions for each $x_1, x_2, ..., x_N$;
- $(a_1) \ a(x) \ge a_0 \text{ and } b(x) \ge b_0 \text{ for some } a_0, b_0 > 0$
- (a_2) $\lambda(x) \geq 0$ for all $x \in \mathbb{R}^N$ and $\lambda(x) > 0$ for all $x \in \Omega$, for some $\Omega \subset \mathbb{R}^N$ such that $|\Omega| < +\infty$.
- $(g_0) g, h \in C^1(\mathbb{R}, \mathbb{R});$
- $(g_1) |g(t)| \le C (1 + |t|^{p-1}), \quad |h(t)| \le C (1 + |t|^{p-1}), \text{ for all } t \in \mathbb{R} \text{ for some } C > 0 \text{ and } p \in (4, 2.2^*)$
- $(g_2) \lim_{t \to 0} \frac{g(t)}{t} = 0, \quad \lim_{t \to 0} \frac{h(t)}{t} = 0;$
- $(g_3) \lim_{|t|\to+\infty} \frac{g(t)}{t^3} = +\infty, \quad \lim_{|t|\to+\infty} \frac{h(t)}{t^3} = +\infty;$
- (g_4) The functions $t \to \frac{g(t)}{t^3}$, $t \to \frac{h(t)}{t^3}$ are strictly increasing in |t|.
- (g_5) There holds $0 \leq G(t) \leq G(|t|)$ and $0 \leq H(t) \leq H(|t|)$, for all $t \in \mathbb{R}$.

Theorem 1.1. Suppose that (a_0) - (a_2) and (g_0) - (g_5) hold. Then, there exists $\theta_0 > 0$ such that System (S_θ) has at least one positive ground state solution, for all $\theta \ge \theta_0$.

References

- A. Ambrosetti, C. Colorado, Bound and ground states of coupled nonlinear Schrödinger equations. C. R. Math. Acad. Sci. Paris 342 (2006), no. 7, 453-458.
- [2] M. Colin, L. Jeanjean, Solutions for a quasilinear Schrödinger equation: a dual approach, Nonlinear Anal. 56 (2004), 213–226. 1
- [3] J. Liu, Z. Q. Wang, X. Wu, Multibump solutions for quasilinear elliptic equations with critical growth. J. Math. Phys. 54 (2013), 121–131.
- [4] J.Q. Liu, Z.Q. Wang, Soliton solutions for quasilinear Schrödinger equations I, Proc. Amer. Math. Soc. 131 (2002), 441–448. 1
- [5] J.Q. Liu, Y.Q. Wang and Z.Q. Wang, Soliton solutions for quasilinear Schrödinger equations II, J. Differential Equations 187 (2003), 473–493.
- [6] J.Q. Liu, Y.Q. Wang and Z.Q. Wang, Solutions for quasilinear Schrödinger equations via the Nehari method, Comm. Partial Differential Equations 29 (2004), 879–901. 1
- [7] L. A. Maia, E. Montefusco, B. Pellacci, Weakly coupled nonlinear Schrödinger systems: the saturation effect, Calc. Var. Partial Differential Equations 46 (2013), no. 1-2, 325-351.
- [8] L. A. Maia, E. Montefusco, B. Pellacci, Positive solutions for a weakly coupled nonlinear Schrödinger system. J. Differential Equations 229 (2006), no. 2, 743-767.
- [9] P.H. Rabinowitz, On a class of nonlinear Schrödinger equations, Z. Angew. Math. Phys. 43 (1992), 270-291.
- [10] E.A.B. Silva, G.F. Vieira, Quasilinear asymptotically periodic Schrödinger equations with subcritical growth, Nonlinear Analysis 72 (2010) 2935–2949.
- [11] M. Yang, Existence of solutions for a quasilinear Schrödinger equation with subcritical nonlinearities, Nonlinear Anal. 75 (2012), 5362–5373.

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